

Generalized Green's Functions.

Consider the nonhomogeneous problem

$$\begin{cases} Au = f(x), & a < x < b & (1) \\ B_1 u(a) = 0, & B_2 u(b) = 0 & (2) \end{cases}$$

when $\lambda = 0$ is an eigenvalue.

We want to find out the solution in terms of a "generalized Green's function".

We have seen that (1) and (2) has ^{no solutions or} infinitely many solutions. Let's assume that $\phi_h(x)$ is an eigenfunction corresponding to $\lambda = 0$, or equivalently that $\phi_h(x)$ is a nontrivial solution of the corresponding homogeneous problem.

To have infinitely many solutions

$$\int_a^b f(x) \phi_h(x) dx = 0$$

$$u = \sum_1^{\infty} u_n \phi_n$$

Recall that

$$u_n \lambda_n = f_n$$

If $\lambda_n = 0 \Rightarrow f_n = 0$ for solutions to exist.

We will illustrate the process of construction of a Generalized Green's function and then ^{construction of} solutions of (1) and (2) through the following example:

$$\begin{cases} \frac{d^2 u}{dx^2} = f(x) = x - \frac{L}{2}, & 0 < x < L. \\ \frac{du}{dx}(0) = 0, & \frac{du}{dx}(L) = 0 \end{cases}$$

Obviously, there are nontrivial solutions for the corresponding homogeneous problem: $\phi_h(x) \equiv 1$. and any multiple

Therefore, for a solution to exist

$$\int_0^L f(x) \phi_h(x) dx = \int_0^L f(x) dx = 0$$

$$\text{If } f(x) = x - L/2 \Rightarrow \int_0^L (x - \frac{L}{2}) dx = \left. \frac{x^2}{2} - \frac{L}{2}x \right|_0^L = 0 \checkmark$$

We observe that the Green's fn. for the BVP given does not exist. In fact, for

$$\text{(for any } x_0) \quad \frac{d^2 g}{dx^2} = \delta(x - x_0), \quad g'(0) = 0, \quad g'(L) = 0 \quad (*)$$

$$\int_0^L \delta(x - x_0) \phi_h(x) dx = \phi_h(x_0) \equiv 1 \neq 0$$

So $\delta(x - x_0) \not\perp \phi_h(x) \equiv 1 \Rightarrow$ There is no soln. (Green's fn.) for (*).

However, we can redefine rhs of (*)
to achieve orthogonality as follows:

$$\frac{d^2 G_m}{dx^2} = \delta(x-x_0) + C \phi_h(x)$$

C is determined from the orthogonality condition

$$0 = \int_0^L \phi_h(x) [\delta(x-x_0) + C \phi_h(x)] dx =$$

$$\Rightarrow 0 = \phi_h(x_0) + C \int_0^L \phi_h^2(x) dx$$

$$\Rightarrow C = \frac{-\phi_h(x_0)}{\int_0^L \phi_h^2(x) dx} \neq 0, \text{ because } \phi_h(x) \neq 0 \text{ nontrivial Soln.}$$

and the related BVP. (nonhomog.)

$$\left\{ \begin{array}{l} \frac{d^2 G_m}{dx^2} = \delta(x-x_0) - \left(\frac{\phi_h(x_0)}{\int_0^L \phi_h^2(x) dx} \right) \phi_h(x) \\ G_m'(0) = 0, \quad G_m'(L) = 0 \end{array} \right.$$

has nontrivial solutions, because rhs $\perp \phi_h(x)$.

In our case, $\phi_h(x) \equiv 1$

$$\Rightarrow C = -\frac{1}{L}$$

related BVP.
$$\begin{cases} \frac{d^2 G_m}{dx^2} = \delta(x-x_0) - \frac{1}{L} \\ G_m'(0) = 0, \quad G_m'(L) = 0 \end{cases}$$

If $x \neq x_0$
$$\frac{d^2 G_m}{dx^2} = -\frac{1}{L}$$

by integrating
$$\frac{dG_m}{dx} = -\frac{1}{L}x + d_{l,r}$$

To satisfy BC at the left.

$$0 = G_m'(0) = d_l \Rightarrow \boxed{d_l = 0}$$

To satisfy BC at the right

$$0 = G_m'(L) = -\frac{L}{L} + d_r \Rightarrow \boxed{d_r = 1}$$

$$\Rightarrow G_m'(x) = \begin{cases} -\frac{x}{L}, & x < x_0 \\ -\frac{x}{L} + 1, & x > x_0 \end{cases} \quad (**)$$

The jump condition in the derivative at $x = x_0$

$$\frac{dG_m}{dx}(x_0+\epsilon) - \frac{dG_m}{dx}(x_0-\epsilon) = -\frac{1}{L} [x_0+\epsilon - x_0+\epsilon] + 1$$

$$\Rightarrow \epsilon \rightarrow 0 \quad \boxed{G_m'(x_0^+) - G_m'(x_0^-) = 1}$$

It's automatically satisfied.

Integrating again (**) $\int dx$

$$G_m(x, x_0) = \begin{cases} -\frac{x^2}{2} \frac{1}{L} + h_1(x_0) & x < x_0 \\ -\frac{x^2}{2} \frac{1}{L} + x + h_2(x_0) & x > x_0 \end{cases}$$

Cont. \Rightarrow at $x_0 \Rightarrow G_m(x_0^+, x_0) = G_m(x_0^-, x_0)$.

$$\Rightarrow -\frac{x_0^{+2}}{2} \frac{1}{L} + x_0 + h_2(x_0) = -\frac{x_0^{-2}}{2} \frac{1}{L} + h_1(x_0)$$

$$\Rightarrow \boxed{h_1(x_0) - h_2(x_0) = x_0} \Rightarrow h_1(x_0) = x_0 + h_2(x_0)$$

or

$$G_m(x, x_0) = \begin{cases} -\frac{x^2}{2} \frac{1}{L} + x_0 + h_2(x_0) & x < x_0 \\ -\frac{x^2}{2} \frac{1}{L} + x + h_2(x_0) & x > x_0 \end{cases}$$

$h_2 \rightarrow h_2(x_0)$

by requiring symmetry: $G_m(x, x_0) = G_m(x_0, x)$, $x < x_0$

$$\Rightarrow h_2(x_0) = -\frac{1}{L} \frac{x_0^2}{2} + \beta$$

$$\Rightarrow G_m(x, x_0) = \begin{cases} -\frac{1}{L} \frac{x^2 + x_0^2}{2} + x_0 + \beta, & x < x_0 \\ -\frac{1}{L} \frac{x^2 + x_0^2}{2} + x + \beta, & x > x_0 \end{cases}$$

$$\left. \begin{array}{l} h_2(x_0) + \frac{x_0}{2} \frac{1}{L} = \\ = h_2(x) + \frac{x}{2} \frac{1}{L} = \beta \end{array} \right\}$$

Now, a soln. (among infinitely many) for the BVP given is

$$u(x) = \int_0^L f(x_0) G_m(x, x_0) dx_0$$

In fact, using Green's formula:

$$\int_a^b [uLv - vLu] dx = \left[p \left(u \frac{dv}{dx} - v \frac{du}{dx} \right) \right]_a^b$$

homog BC's

$$\int_a^b \left\{ u(x) \left[\delta(x-x_0) - \frac{\phi_h(x) \phi_h(x_0)}{\int_a^b \phi_h^2(x_0) dx_0} \right] - G_m(x, x_0) f(x) \right\} dx = 0.$$

$L[G_m(x, x_0)]$

\downarrow
 Lu

$$\Rightarrow u(x_0) = \frac{\phi_h(x_0)}{\int_a^b \phi_h^2(x_0) dx_0} \int_a^b \phi_h(x) u(x) dx + \int_a^b G_m(x, x_0) f(x) dx$$

Reversing Variables $x_0 \leftrightarrow x$. and using symmetry of $G_m(x, x_0)$

$$u(x) = \frac{\phi_h(x)}{\int_a^b \phi_h^2(\bar{x}) d\bar{x}} \int_a^b \phi_h(x_0) u(x_0) dx_0 + \int_a^b G_m(x, x_0) f(x_0) dx_0$$

multiple of homog. soln.

So particular soln.

$$u(x) = \int_a^b G_m(x, x_0) f(x_0) dx_0$$